PHY 554. Homework 1.

Handed: August 27 Return by: September 10
Bring solution to class or -better- email solutions to <u>vladimir.litvinenko@stonybrook.edu</u>

HW 1.1 (3 points): Find available energy (so called C.M. energy) for a head-on collision in these scenarios:

- (a) In CEBAF, polarized 12 GeV electrons collide with protons at rest;
- (b) In one of scenarios for future collides, CERN consider collide 50 TeV protons with 180 GeV electrons;

The rest energy of proton is 938.257 MeV, and rest energy of electron is 0.511 MeV. **Solution.**

First, the remined of the center of mass energy available for creating new particle:

$$p^{\mu} = \left(p_{o} = \frac{E}{c}, \vec{p}\right); p^{\mu}p_{\mu} = p_{o}^{2} - \vec{p}^{2} = p_{cm}^{2} = \left(M_{cm}c\right)^{2};$$

$$E = E_{1} + E_{2} = \gamma_{1}m_{1}c^{2} + \gamma_{2}m_{2}c^{2}; \vec{p} = \vec{p}_{1} + \vec{p}_{2} = \gamma_{1}m_{1}\vec{v}_{1} + \gamma_{2}m_{2}\vec{v}_{2}$$

$$\left(M_{cm}c^{2}\right)^{2} = \left(E_{1} + E_{2}\right)^{2} - \left(\vec{p}_{1} + \vec{p}_{2}\right)^{2}c^{2} =$$

$$\left(E_{1}^{2} - \vec{p}_{1}^{2}c^{2}\right) + \left(E_{2}^{2} - \vec{p}_{2}^{2}c^{2}\right) + 2\left(E_{1}E_{2} - 2\vec{p}_{1}\vec{p}c^{2}\right)$$

$$\left(\frac{2}{3}c^{2} + \frac{2}{3}c^{2}c^{2}\right) + \frac{2}{3}c^{2}c^{2} + \frac{2}{3}c^{2}c^{2}c^{2} + 2\left(E_{1}E_{2} - 2\vec{p}_{1}\vec{p}c^{2}\right)$$

and using $E_{1,2}^2 - \vec{p}_{1,2}^2 c^2 = (m_{1,2}^2 c^2)^2$

we get

$$\left(M_{cm}c^2 \right)^2 = \left(m_1c^2 \right)^2 + \left(m_2c^2 \right)^2 + 2 \left(\gamma_1 m_1c^2 \right) \left(\gamma_2 m_2c^2 \right) \left(1 - \vec{\beta}_1 \vec{\beta}_2 \right); \ \vec{\beta}_{1,2} = \frac{\vec{v}_{1,2}}{c}$$

$$M_{cm}c^2 = \sqrt{\left(m_1c^2 \right)^2 + \left(m_2c^2 \right)^2 + 2E_1E_2 \left(1 - \vec{\beta}_1 \vec{\beta}_2 \right)}$$

The available c.m. energy depends not only on particles energy but dimensionless term $\vec{\beta}_1 \vec{\beta}_2$. The product is equal to zero is one of the particles is at rest or particles collide at 90 degrees and available energy is maximized for head-on collisions when

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$$\vec{\beta}_1 \vec{\beta}_2 = - \left| \vec{\beta}_1 \right| \left| \vec{\beta}_2 \right| = -\beta_1 \beta_2 \Rightarrow M_{cm} c^2 = \sqrt{\left(m_1 c^2 \right)^2 + \left(m_2 c^2 \right)^2 + 2E_1 E_2 \left(1 + \beta_1 \beta_2 \right)};$$

(a) Electrons are moving with velocity close to speed of light, but protons are at rest (β_2 =0) and E₂=m₂c² and we get

$$M_{cm}c^{2} = \sqrt{(m_{1}c^{2})^{2} + (m_{2}c^{2})^{2} + 2E_{1}m_{2}c^{2}} = \sqrt{(0.000511)^{2} + (2 \cdot 12 + 0.938257) \cdot 0.938257} = 4.837 \text{ GeV}$$

If we use an approximate formular neglecting the c.m. of the proton, we would miss ~ 90 keV:

$$M_{cm}c^2 \cong \sqrt{2E_1 m_2 c^2} = 4.745 \ GeV$$

(b) Beams are colliding head-on, i.e. we need to use

$$M_{cm}c^2 = \sqrt{(m_1c^2)^2 + (m_2c^2)^2 + 2E_1E_2(1 + \beta_1\beta_2)}$$

Both beams are ultra-relativistic with $\gamma_e=3.52\times10^5$ and $\gamma_p=5.33\times10^4$. Off cause, we can calculate β for both beams and find that $\beta_e=1-4.03\times10^{-12}$ and $\beta_e=1-1.76\times10^{-10}$ to get $M_{cm}c^2=\sqrt{\left(m_1c^2\right)^2+\left(m_2c^2\right)^2+2E_1E_2\left(1+\beta_1\beta_2\right)}=6.00000000731~TeV$

$$M_{cm}c^2 = \sqrt{(m_1c^2)^2 + (m_2c^2)^2 + 2E_1E_2(1 + \beta_1\beta_2)} = 6.0000000731 \ TeV$$

But in this case we really can neglect first two terms and replace βs by 1s:

$$M_{cm}c^2 \cong 2\sqrt{E_1 E_2} = 6 \text{ TeV}$$

with very tiny $1.2x10^{-8}$ relative difference from exact answer of 73 keV.

HW 1.2 (2 points): Let's consider a futuristic case that humanity decided – and figured out how - to build a storage ring surrounding the equator, i.e. with circumference of 40,075 km.

- 1 point: If bending dipole magnets fill 75% of the ring circumference, and we used dipole magnetic field of 8.3 T, demonstrated in LHC, what would be energy of the circulating protons?
- (b) 1 point: What magnetic field would be needed if storage ring with the same energy of the ring should fit at Earth largest continent of Asia. Asia could fit a ring with radius of 2,500 km. Assume 75% filling factor.

P.S. USA can fit ring with radius of 1.465 km.

Solution:

- (a) Circumference of 40.075 km means that average radius is R=6.378 km and ff=75% of the trajectory has radius of ρ =4,783.6 km. With B_o=8.3 T it would correspond to B ρ =39.704 T km and pc=11.90 PeV (1 PeV = 10^{15} eV)– not to shabby at all. Beam would ultra-relativistic and E= pc=11.90 PeV with very high accuracy/
- (b) To have the same energy, i.e. the same momentum and beam rigidity $B\rho=39.704$ T km. With the same filling factor ρ scale proportionally to the average radius of the storage ring:

$$B_{req} = \frac{B\rho}{\rho_{asia}} = \frac{B_o \cdot \rho_{earth}}{\rho_{asia}} = \frac{B_o \cdot ff \cdot R_{earth}}{ff \cdot R_{asia}} = B_o \frac{R_{earth}}{R_{asia}} = 21.18 T$$

It would require field exceeding 36 T to fit it in the USA, or operate it at 2.7 PeV suing 8.3T magnetic field.

HW 1.3 (2 points): For a classical microtron with orbit factor k=1 and energy gain per pass of 1.022 MeV and operational RF frequency 3 GHz (3 x 10⁹ Hz) find required magnetic field. What will be radius of first orbit in this microtron?

Hint: Note that rest energy of electron with $\gamma=1$ is 0.511 MeV. This is energy gain per pass will define available n numbers in eq. (2.6)

Solution: The main equations for the microtrons are derived from following equations for beam going on nth pass through the cavity:

$$\begin{split} E_n = mc^2 + n \cdot e \cdot V_{rf} &\Rightarrow \gamma_n = \frac{E_n}{mc^2} = 1 + n \cdot \frac{e \cdot V_{rf}}{mc^2}; \, \beta_n = \frac{\mathbf{v}_n}{c} = \sqrt{1 - \gamma_n^{-2}}; \\ p_n c = \beta_n \cdot E_n = e \cdot B \cdot \rho_n; \, \rho_n = \frac{\beta_n \cdot E_n}{e \cdot B}; \, T_n = \frac{2\pi \cdot \rho_n}{\mathbf{v}_n} = \frac{2\pi \cdot E_n}{e \cdot B} \cdot \frac{\beta_n}{\mathbf{v}_n} = \frac{2\pi \cdot E_n}{e \cdot B \cdot c} = \frac{2\pi \cdot mc}{e \cdot B} \cdot \left(1 + n \cdot \frac{e \cdot V_{rf}}{mc^2}\right). \end{split}$$

where T_n is the time required for particle turn around nt^h path. The resonant condition for microtron is that each pass takes integer number of RF cycles:

$$\begin{split} T_n &= \frac{2\pi \cdot mc}{e \cdot B} \left(1 + n \cdot \frac{e \cdot V_{rf}}{mc^2} \right) = N_n \cdot T_{RF} = N_n \cdot \frac{1}{f_{RF}}; \\ N_n &= \frac{2\pi \cdot mc \cdot f_{RF}}{e \cdot B} \left(1 + n \cdot \frac{e \cdot V_{rf}}{mc^2} \right) = \frac{2\pi \cdot mc^2}{e \cdot B \cdot \lambda_{RF}} \cdot \left(1 + n \cdot \frac{e \cdot V_{rf}}{mc^2} \right); \ \lambda_{RF} = \frac{c}{f_{RF}}; \\ N_n &= A \cdot C_n; \ A = \frac{2\pi \cdot mc^2}{e \cdot B \cdot \lambda_{RF}}; \ C_n = 1 + n \cdot \frac{e \cdot V_{rf}}{mc^2}. \end{split}$$

For product of two numbers A and C_n to be integers for all n=1,2,... it is required that A is an integer and that

$$A \cdot \frac{e \cdot V_{rf}}{mc^2}$$

is also integer. In our case, $eV_{rf}/mc^2=2$ and second condition is satisfied if A is an integer, which determining

$$B = \frac{2\pi}{A} \cdot \frac{mc^2}{e \cdot \lambda_{RE}}$$

We are selecting A=1 to get

$$B = \frac{2\pi \cdot mc^2}{e \cdot \lambda_{RE}}; \ 2\pi \cdot \frac{mc^2}{e} = 2\pi \cdot \frac{0.511[MeV]}{0.29979} = 10.71[kGs\ cm];$$

$$\lambda_{RF} = \frac{c}{f_{RF}} = \frac{2.9979 \cdot 10^{10} \left[\frac{cm}{sec} \right]}{3 \cdot 10^9 \left[\frac{1}{sec} \right]} = 9.993 \ cm; \ B = 1.071 \ kGs \ (0.1071 \ T)$$

At first orbit (n=1), the beam energy is E=3 mc² =1.533 MeV and particle rigidity

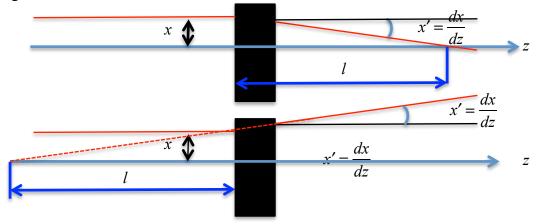
$$pc = \sqrt{E^2 - (mc^2)^2} = 1.445 \text{ MeV}; B\rho[kGs \ cm] = \frac{pc[MeV]}{0.29979} = 4.821 \text{ kGs } cm$$

$$B = 1.071 \ kGs; \ \rho = \frac{B\rho}{B} = 4.498 \ cm$$

HW 1.4 (8 point): Let's first determine an effective focal length, F, of a paraxial (e.g. small angles!) focusing object (a black-box) as ratio between a parallel displacement of trajectory at its entrance to corresponding change of the angle at its exit (see figure below):

$$F = -\frac{x}{x}; x' \equiv \frac{dx}{dz}$$

see figure below for



Let consider a triplet of thin lenses: horizontal focusing (F) and defocusing (D) lenses with centers are separated by distance L as in Fig. 1. Since quadrupole focusing is opposite for horizontal and vertical directions, the lattice - in accelerator lingo it is magnetic structure - looks like (a) in horizontal direction and (b) in vertical direction. Let's assume that that in Fig 1(a) Q_1 are focusing lenses in horizontal plane with focal length of $F=+1/q_1$ and, naturally, they are defocusing in vertical direction, Fig 1 (b), with focal length of $F=-1/q_1$. Similarly, Q_2 is defocusing in horizontal plane (Fig 1(a)) with focal length of $F=-1/q_2$ and is focusing in vertical plane with focal length of $F=+1/q_2$.

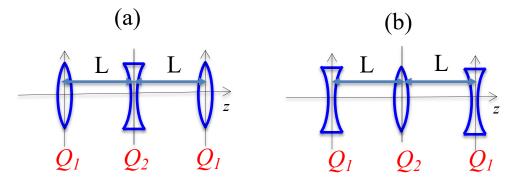
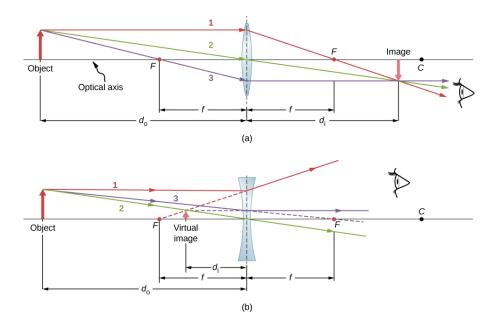


Fig.1. Triplet of short quadrupole lenses: (a) in horizontal direction FDF and (b) in vertical direction DFD.

- 1. (5 points) Find of horizonal (FDF) and vertical (DFD) focal lengths of the triplet.
- 2. (3 points) Find ratio between $\mathbf{q_1}$ and $\mathbf{q_2}$ when horizontal and vertical focal lengths are equal, specifically in a form of $\mathbf{q_2}$ = $\mathbf{f}(\mathbf{q_1},\mathbf{L})$.

P.S. Definition (picture) of thin lens:



Solution. Instead of doing repetitive calculations of effect on x and x' by lenses and drifts (space between lenses) we will write it in compact, matrix form.

The thin lens with focal length F=1/q does not change position but changes angle if trajectory:

$$x_{out} = x_{in}, \ x'_{out} = x'_{in} - \frac{x_{in}}{f} = x'_{in} - q \cdot x_{in} \Rightarrow \begin{bmatrix} x_{out} \\ x'_{out} \end{bmatrix} = \begin{bmatrix} 1 & 0 \\ -q & 1 \end{bmatrix} \cdot \begin{bmatrix} x_{in} \\ x'_{in} \end{bmatrix}$$

Sign of q determines is this focusing (q>0) or defocusing (q<0) length.

The drift of length L does not changes angle of trajectory but changes position:

$$x_{out} = x_{in} + L \cdot x'_{in}, \ x'_{out} = x'_{in} \Rightarrow \begin{bmatrix} x_{out} \\ x'_{out} \end{bmatrix} = \begin{bmatrix} 1 & L \\ 0 & 1 \end{bmatrix} \cdot \begin{bmatrix} x_{in} \\ x'_{in} \end{bmatrix}$$

Hence the triplet will result in multiplying 5 matrices – 3 of thin lenses and two drift between them. For horizontal case (a) we have

$$\begin{bmatrix} x_{out} \\ x'_{out} \end{bmatrix} = M_a \cdot \begin{bmatrix} x_{in} \\ x'_{in} \end{bmatrix} M_a = \begin{bmatrix} 1 & 0 \\ -q_1 & 1 \end{bmatrix} \cdot \begin{bmatrix} 1 & L \\ 0 & 1 \end{bmatrix} \cdot \begin{bmatrix} 1 & 0 \\ q_2 & 1 \end{bmatrix} \cdot \begin{bmatrix} 1 & L \\ 0 & 1 \end{bmatrix} \cdot \begin{bmatrix} 1 & 0 \\ -q_1 & 1 \end{bmatrix} = \begin{bmatrix} 1 - L(2q_1 - q_2) - L^2q_1q_2 & L(2 + Lq_2) \\ -(1 - Lq_1)(q_1(2 + Lq_2) - q_2) & 1 - L(2q_1 - q_2) - L^2q_1q_2 \end{bmatrix}$$

For vertical case (b) we just need to reverse signs of q's:

$$\begin{bmatrix} x_{out} \\ x'_{out} \end{bmatrix} = M_b \cdot \begin{bmatrix} x_{in} \\ x'_{in} \end{bmatrix} M_b = \begin{bmatrix} 1 & 0 \\ q_1 & 1 \end{bmatrix} \cdot \begin{bmatrix} 1 & L \\ 0 & 1 \end{bmatrix} \cdot \begin{bmatrix} 1 & 0 \\ -q_2 & 1 \end{bmatrix} \cdot \begin{bmatrix} 1 & L \\ 0 & 1 \end{bmatrix} \cdot \begin{bmatrix} 1 & 0 \\ q_1 & 1 \end{bmatrix} = \begin{bmatrix} 1 + L(2q_1 - q_2) - L^2q_1q_2 & L(2 - Lq_2) \\ (1 + Lq_1)(q_1(2 - Lq_2) - q_2) & 1 + L(2q_1 - q_2) - L^2q_1q_2 \end{bmatrix}$$

We define the focal length if the "black box" as distance from it exit of the system till it crossing the axis for initial trajectory without angle: x'_{in}=0:

$$\begin{bmatrix} x_{out} \\ x'_{out} \end{bmatrix}_{a} = M_{a} \cdot \begin{bmatrix} x_{o} \\ 0 \end{bmatrix} = \begin{bmatrix} 1 - L(2q_{1} - q_{2}) - L^{2}q_{1}q_{2} \\ -(1 - Lq_{1})(q_{1}(2 + Lq_{2}) - q_{2}) \end{bmatrix} \cdot x_{o};$$

$$\begin{bmatrix} x_{out} \\ x'_{out} \end{bmatrix}_{b} = M_{b} \cdot \begin{bmatrix} x_{o} \\ 0 \end{bmatrix} = \begin{bmatrix} 1 + L(2q_{1} - q_{2}) - L^{2}q_{1}q_{2} \\ (1 + Lq_{1})(q_{1}(2 - Lq_{2}) - q_{2}) \end{bmatrix} \cdot x_{o};$$

$$F_{a} = -\frac{(x_{out})_{a}}{(x'_{out})_{a}} = \frac{1 - L(2q_{1} - q_{2}) - L^{2}q_{1}q_{2}}{(1 - Lq_{1})(q_{1}(2 + Lq_{2}) - q_{2})};$$

$$F_{b} = -\frac{(x_{out})_{b}}{(x'_{out})_{b}} = \frac{1 + L(2q_{1} - q_{2}) - L^{2}q_{1}q_{2}}{-(1 + Lq_{1})(q_{1}(2 - Lq_{2}) - q_{2})}$$

It will be sufficient to write $F_a = F_b$ to get the full points, but if you go and solved it:

$$Solve \left[\frac{1 - L\left(2q_{1} - q_{2}\right) - L^{2}q_{1}q_{2}}{\left(1 - Lq_{1}\right)\left(q_{1}\left(2 + Lq_{2}\right) - q_{2}\right)} = -\frac{1 + L\left(2q_{1} - q_{2}\right) - L^{2}q_{1}q_{2}}{\left(1 + Lq_{1}\right)\left(q_{1}\left(2 - Lq_{2}\right) - q_{2}\right)} \right];$$

$$q_{2} = \frac{1}{2L} \frac{1 - 3u \pm \sqrt{1 + 2u^{2} - 15u^{4} + 16u^{6}}}{u^{3} - u}; u = Lq_{1}$$

I will double your points! It turns out that solution with minus sign in front of the square root is focusing when 0 < u < 0.708 and defocusing at u > 0.708. At u = 0.708, the focal spot is located at the exit of the triplet, i.e. parallel beam merges in a point at its exit. Solution with positive sign is defocusing at 0 < u < 1 and focusing at u > 0.