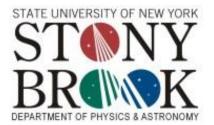
Low Emittance Storage Ring for Light Source

Sukho Kongtawong PHY 554 Fall 2016



Content

- Brightness and emittance
- Radiative effect and emittance
- Theory
 - Theoretical Minimum Emittance (TME) cell
 - Double-bend achromat (DBA)
- Multi-bend achromat (MBA)
 - MAX IV (7BA)
- Upgrading of Synchrotron Storage Rings

Brightness and emittance

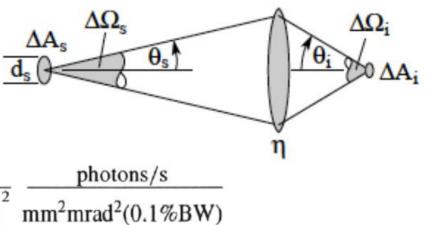
• All beamline experiments require high brightness of the light $B_{\Delta\omega/\omega} = \frac{\Delta P}{\Delta A \Delta \Omega \Delta \omega/\omega},$

 $7.25 \times 10^6 N^2 N^2 I(\Lambda)$

which depends on emittance

• Example: Undulator





$$\bar{B}_{\Delta\omega/\omega}(0) = \frac{1.25 \times 10^{\circ} \gamma N I(R)}{\sigma_x(\text{mm})\sigma_y(\text{mm}) \left(1 + \frac{{\sigma'_x}^2}{\theta_{\text{cen}}^2}\right)^{1/2} \left(1 + \frac{{\sigma'_y}^2}{\theta_{\text{cen}}^2}\right)^{1/2}} \cdot \frac{K^2 f(K)}{(1 + K^2/2)^2} \frac{\text{photoms/s}}{\text{mm}^2 \text{mrad}^2(0.1\%\text{B}^2)}$$
where $\sigma = \sqrt{\epsilon\beta}$ and $\sigma' = \sqrt{\epsilon/\beta}$

v2 civ

Damping term $\alpha_{\chi} = \frac{U_0}{2T_0E}(1-\overline{D}) = \frac{U_0}{2T_0E}J_{\chi}$

Radiative Effects

- Balance between Beam damping and Quantum fluctuation
- Emittance can be determined by

$$\varepsilon_x = C_q \gamma^2 \frac{I_5}{I_2 - I_4} = C_q \gamma^2 \frac{\langle H(s)/\rho^3 \rangle}{J_x \langle 1/\rho^2 \rangle},$$

where

$$H(s)=\gamma D^2+2\alpha DD'+\beta D'^2,$$

which can be adjusted by vary magnetic strength

$$\varepsilon_x = C_q \gamma^2 \frac{\langle H(s)/\rho^3 \rangle}{J_x \langle 1/\rho^2 \rangle},$$

Radiative Effects

- Most Storage rings use identical bending radius
- Strongly focusing machines generally have $J_{\chi} \approx 1$
- The emittance becomes

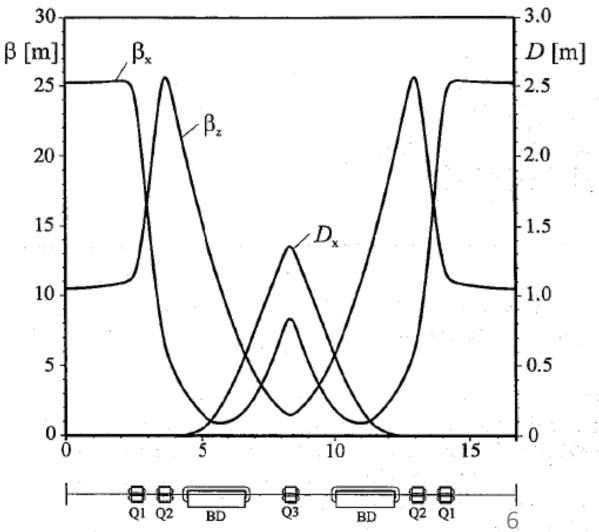
$$\varepsilon_{\chi} = C_q \frac{\gamma^2}{\rho l} \int_{0}^{l} H(s) \, ds$$

where l is length of bending magnets

• Optimize H(s) to reduce the emittance

Chaseman-Green Lattice DBA

- Designed by Renate Chaseman and George K. Green of BNL in mid-1970s
- Aka "double bend achromat"
- Dispersion, D = 0, D' = 0 in straight section for IDs
- Optimize Betatron function



Klaus Wille, "The Physics of Particle Accelerators: An Introduction", Oxford University press

Evolution of optical parameters in Bending

• Find evolution of D(s), D'(s), $\beta(s)$, $\alpha(s)$ for

$$H(s) = \gamma D^2 + 2\alpha D D' + \beta D'^2$$

• For dispersion

$$\begin{pmatrix} D(s) \\ D'(s) \\ 1 \end{pmatrix} = \begin{pmatrix} \cos(s/\rho) & \rho \sin(s/\rho) & \rho(1 - \cos(s/\rho)) \\ -\sin(s/\rho) / \rho & \cos(s/\rho) & \sin(s/\rho) \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} D_0 \\ D'_0 \\ 1 \end{pmatrix}$$

and $D_0 = 0, D'_0 = 0$ gives
$$D(s) = \rho(1 - \cos(s/\rho)) \approx \frac{s^2}{2\rho}$$
$$D'(s) = \sin(s/\rho) \approx \frac{s}{\rho}$$

Klaus Wille, "The Physics of Particle Accelerators: An Introduction", Oxford University press

Evolution of optical parameters in Bending

For Betatron function (transformation Matrix for drift space)

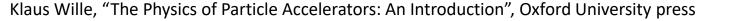
$$\begin{pmatrix} \beta(s) & -\alpha(s) \\ -\alpha(s) & \gamma(s) \end{pmatrix} = \begin{pmatrix} 1 & s \\ 0 & 1 \end{pmatrix} \begin{pmatrix} \beta_0 & -\alpha_0 \\ -\alpha_0 & \gamma_0 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ s & 1 \end{pmatrix}$$

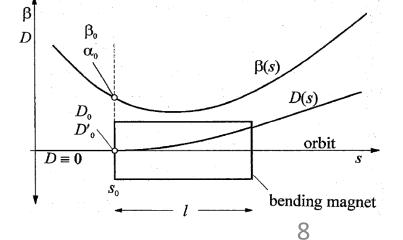
gives

$$\beta(s) = \beta_0 - 2\alpha_0 s + \gamma_0 s^2$$

$$\alpha(s) = \alpha_0 + \gamma_0 s$$

$$\gamma(s) = \gamma_0$$





Emittance

Hence

$$H(s) = \frac{1}{\rho^2} \left(\frac{\gamma_0}{4} s^4 - \alpha_0 s^3 + \beta_0 s^2 \right)$$
$$\varepsilon_x = C_q \frac{\gamma^2}{\rho l} \int_0^l H(s) \, ds$$

Finally,

$$\varepsilon_{x} = C_{q} \gamma^{2} \Theta^{3} \left(\frac{\gamma_{0} l}{20} - \frac{\alpha_{0}}{4} + \frac{\beta_{0}}{3l} \right)$$

where $\Theta = l/\rho$

Minimum Emittance

Differentiate respects to initial condition

$$\frac{\partial \varepsilon_x}{\partial \beta_0} = 0 \text{ and } \frac{\partial \varepsilon_x}{\partial \alpha_0} = 0$$

give

So

$$\frac{1+\alpha_0^2}{\beta_0^2}\frac{l^2}{20} + \frac{1}{3} = 0 \text{ and } \frac{\alpha_0}{\beta_0}\frac{l}{10} - \frac{1}{4} = \frac{1}{2}$$
$$\beta_0 = \sqrt{\frac{12}{5}}l, \ \alpha_0 = \sqrt{15}$$

Klaus Wille, "The Physics of Particle Accelerators: An Introduction", Oxford University press

0

Minimum Emittance

For Chasman-Green lattice (double bend achromat)

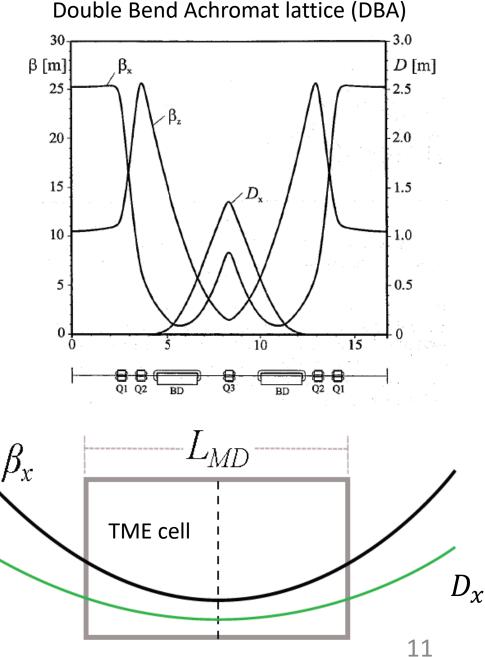
 $\varepsilon_x \text{ [nm-rad]} = 94.7 \ \Theta^3 E^2 \text{[GeV]}$

Note that if it's not Achromat lattice when $D \neq 0$ and $D' \neq 0$

$$\varepsilon_x \text{ [nm-rad]} = \frac{31.65}{J_x} \Theta^3 E^2 \text{[GeV]}$$

This is the theoretical minimum emittance (TME)

Dieter Einfeld (2014), "Multi-bend Achromat Lattices for Storage Ring Light Sources", Synchrotron Radiation News



Multi Bend Achromat (MBA)

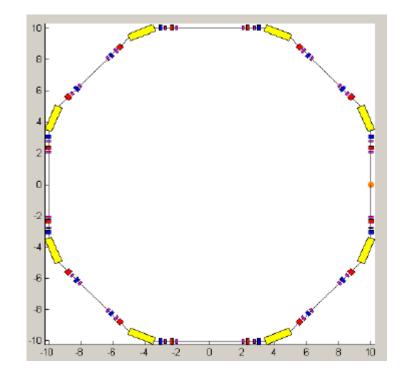
- D. Einfeld *et al*. proposed the use of MBA in 1993.
- From

$$\varepsilon_x = C_q \frac{\gamma^2}{J_x} \Theta^3 F,$$

where F is factor depends on Betatron function

- $\Theta = 2\pi/N$, N is number of bending
- Hence

$$\varepsilon_{\chi} \propto \frac{E^2}{N^3}$$



• Increase $N \rightarrow \text{reduce } \Theta \rightarrow \text{reduce } \varepsilon_{\chi}$

MBA

2.5

2

Emittance (nmrad)

0.5

0 1.2

1.3

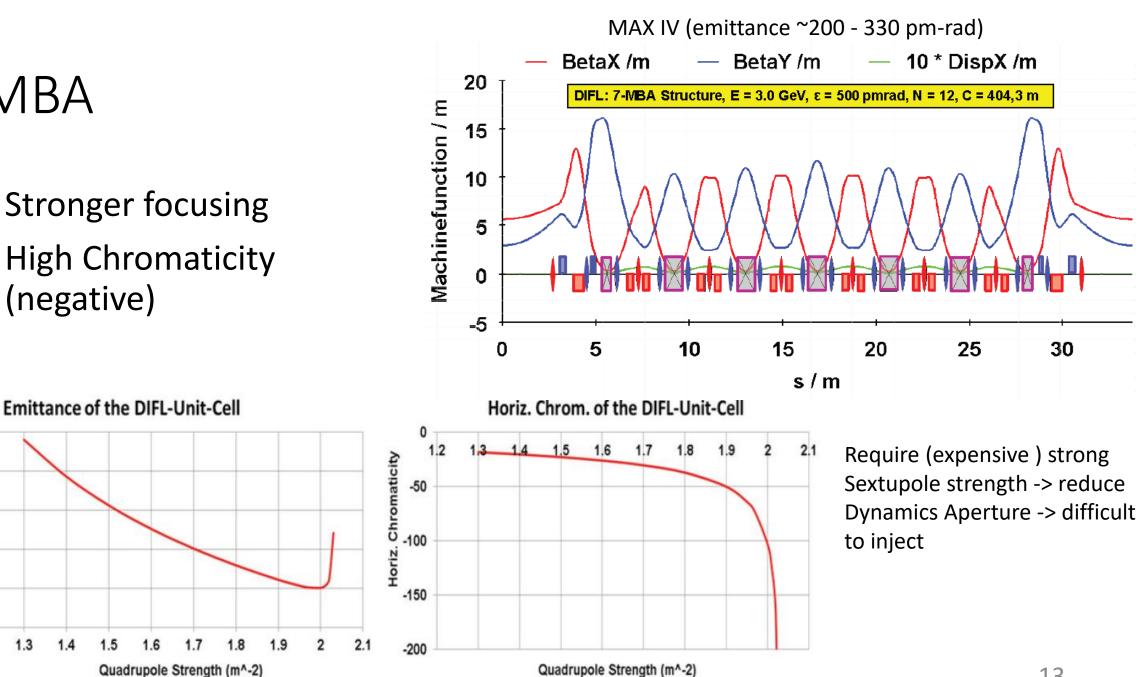
1.4

1.5

1.6

1.7

- Stronger focusing
- High Chromaticity (negative)



Dieter Einfeld (2014), "Multi-bend Achromat Lattices for Storage Ring Light Sources", Synchrotron Radiation News

MBA

- More bending -> reduce space for IDs
- Require combine magnet -> increase J_x which reduce ε_x

$$\alpha_{E} = \frac{U_{0}}{2T_{0}E} (2 + \overline{D}) = \frac{U_{0}}{2T_{0}E} J_{E}$$

$$\alpha_x = \frac{U_0}{2T_0 E} (1 - \overline{D}) = \frac{U_0}{2T_0 E} J_x$$

Hence $J_x + J_E = 3$ constant

• Increase J_{χ} will reduce J_E -> increase energy spread

MAX IV

528 m circumference

The world's first MBA storage ring, Lund, Sweden

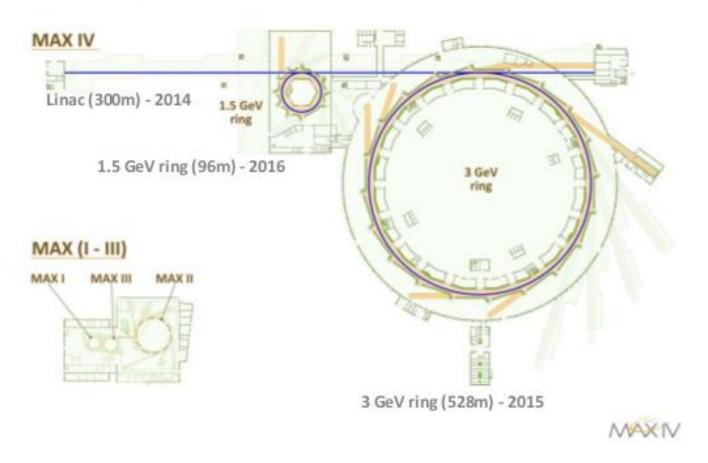


MAX IV

Consists of

- 3 GeV LINAC (300 m)
 - 10 Hz for full energy injection
 - XFEL
- 1.5 GeV 96 m circumference
 - Soft x-ray and UV users
- 3 GeV: 528 m circumference
 - Hard x-ray users

The layout of MAX IV



MAX IV

Comparison between

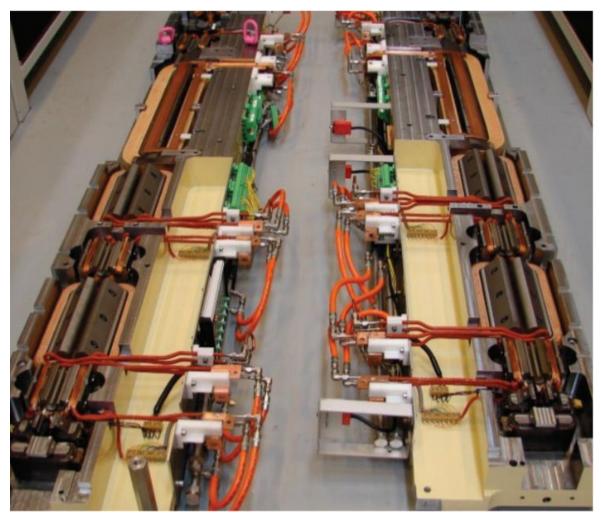
- NSLS-II
- MAX IV



Parameters	NSLS-II	MAX IV
Energy (GeV)	3	3
Circumference (m)	792	528
Current (mA)	500	500
No. of Straight Section	15 (9.3 m) / 15 (6.6 m)	20
Horizontal Emittance (nm rad)	0.55	~0.2 - 0.33
Vertical Emittance (pm)	8	2 - 8

MAX IV: Magnets

- In order to reach low emittance with small circumference
 - Compact magnet
 - Reduce magnets gap -> Strong magnetic field
 - Shorter
- Integrated units
 - BM, QM pole are machined out of a pair of iron blocks
 - Each unit holding all the magnets of a complete cell



Top and bottom parts of multiple function magnet of MAX IV Credit: Danfysik

Magnets

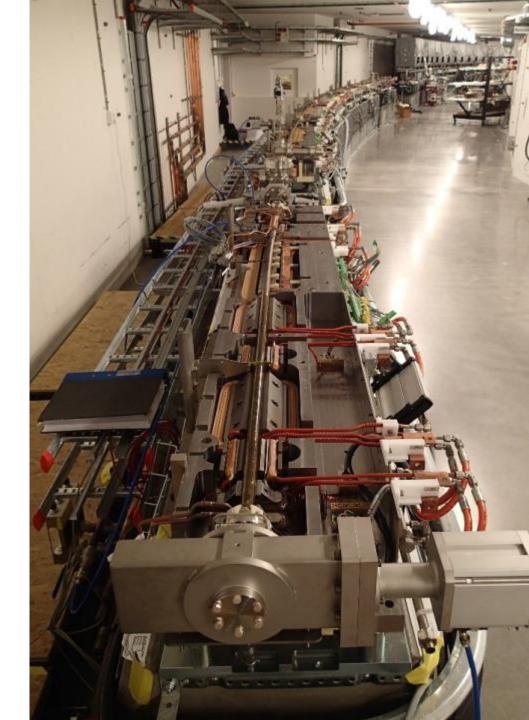


Separated function magnets (NSLS-II)

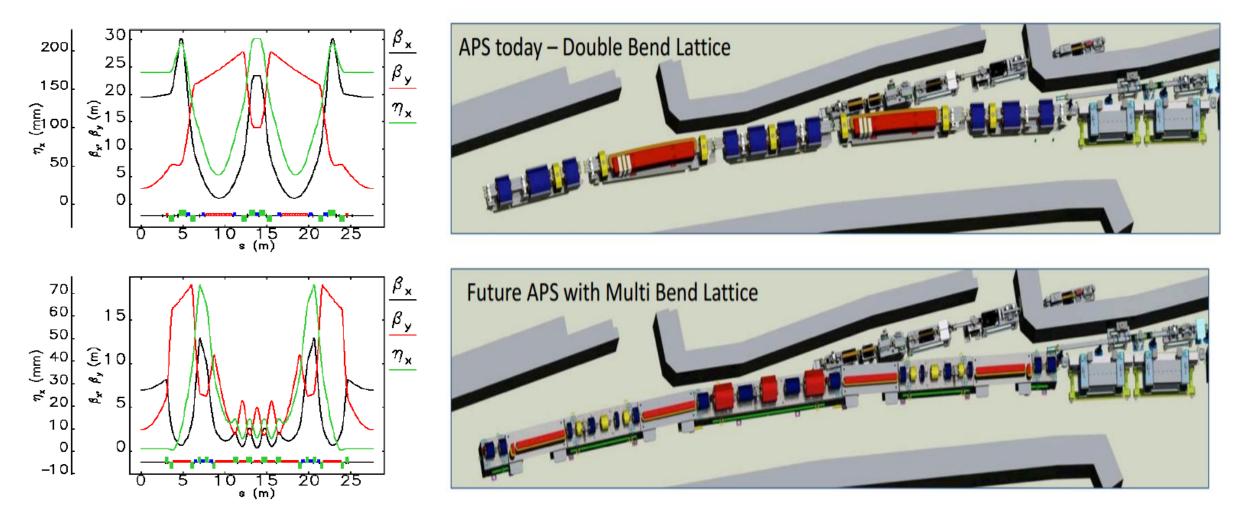
Two complete 7-bend achromats at MAX IV

MAX IV: Chamber

- Compact magnet design lead to narrow low-conductance vacuum chambers
 - Distributed pumping and heat load
- Heat load problem
 - Use copper as a chamber
 - Water cooling
- Pumping problem
 - Provide by non-evaporable getter (NEG) coating at the inner surface
 - Reduce number of pumps and absorbers



Advance Photon Source (APS) upgrading



Stuart Henderson, "APS Upgrade Overview and Status", APS Users Organization/Partner User Council Joint Meeting July 9, 2014

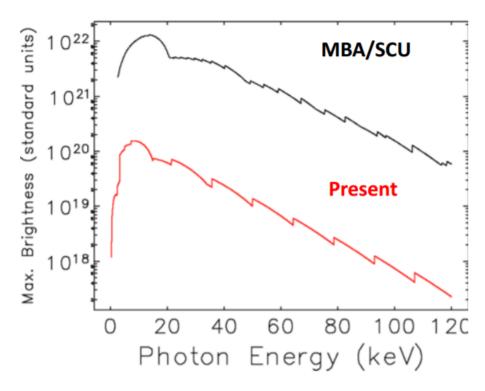
Advance Photon Source (APS) upgrading

- 1104 m circumference
- Upgrade
 - APS-U MBA
 - 67 pm-rad
- 100 fold increase in brightness and coherent flux

Radiation-integral-related quantities			
Beam energy	7	6	GeV
Natural emittance	2527.5	66.9	$_{\rm pm}$
Energy spread	0.095	0.096	%

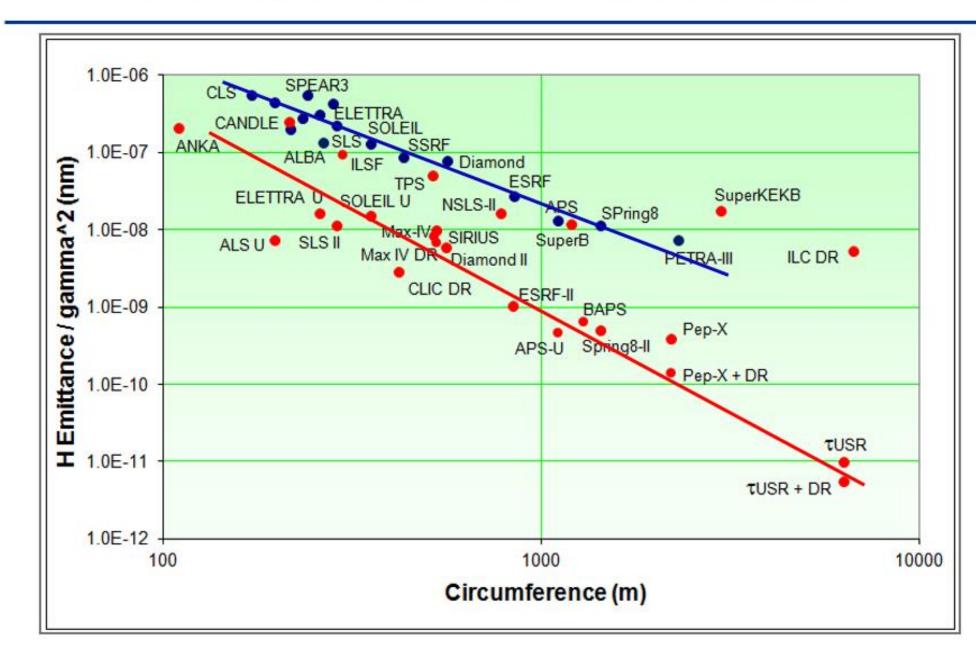
DBA

MBA



Stuart Henderson, "APS Upgrade Overview and Status", APS Users Organization/Partner User Council Joint Meeting July 9, 2014

Survey of low emittance lattices



Credit: "Status of the Diamond Light Source upgrade", R. Bartolini, et al., EuCARD2 topical workshop Barcelona, 23 April 2015

Summary

- Theoretical Minimum Emittance (TME) consider only on emittance
 - Difficult to reach
- In reality, other parameters need to be considered
 - Non-linear term e.g. dynamics aperture
 - Space for IDs
 - Costs e.g. high-tech magnets
- MBA can reduce emittance because reduce the bending angle
- Many Synchrotrons around the world plan to upgrade to MBA

Reference

- PHY 554 CASE course, Stony Brook University, Fall 2016
- Klaus Wille, "The Physics of Particle Accelerators: An Introduction", Oxford University press
- Dieter Einfeld (2014), "Multi-bend Achromat Lattices for Storage Ring Light Sources", Synchrotron Radiation News, 27:6, 4-7, DOI: 10.1080/08940886.2014.970929
- Pedro F. Tavares, "The MAX IV storage ring project", Journal of Synchrotron Radiation ISSN, 1600-5775
- Thapakron Pulampong, "Ultra-low Emittance Lattice Design for Advanced Synchrotron Light Sources", thesis submitted at the University of Oxford
- R. Bartolini. A. Alekou, M. Apollonio, et al., "Status of the Diamond Light Source upgrade", EuCARD2 topical workshop Barcelona, 23 April 2015
- Stuart Henderson, "APS Upgrade Overview and Status", APS Users Organization/Partner User Council Joint Meeting July 9, 2014

Extra Slide: Chromaticity

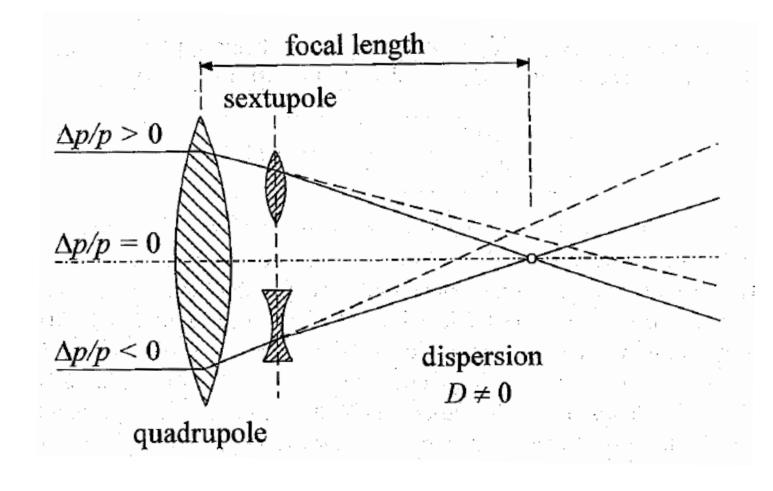
Chromaticity leads to tune shifted

$$\xi \equiv \frac{\Delta Q}{\Delta p/p} = \oint k(s)\beta(s)ds$$

- Sextupole magnets
 - Compensate the Chromaticity

$$k_{\text{sext}} = m D \frac{\Delta p}{p}$$

Where m is sextupole strength



Extra Slide: NEG coating

- Non evaporable getters (NEG), based on the principle of metallic surface sorption of gas molecules, are mostly porous alloys or powder mixtures of Al, Zr, Ti, V and Fe. They help to establish and maintain vacuums by soaking up or bonding to gas molecules that remain within a partial vacuum.
- They are important tools for improving the performance of many vacuum systems.
- the NEG coating can be applied even to spaces that are narrow and hard to pump out, which makes it very popular in particle accelerators where this is an issue.
- The NEG acts as a getter or getter pump that is able to reduce the pressure to less than 10⁻⁷ Pa.