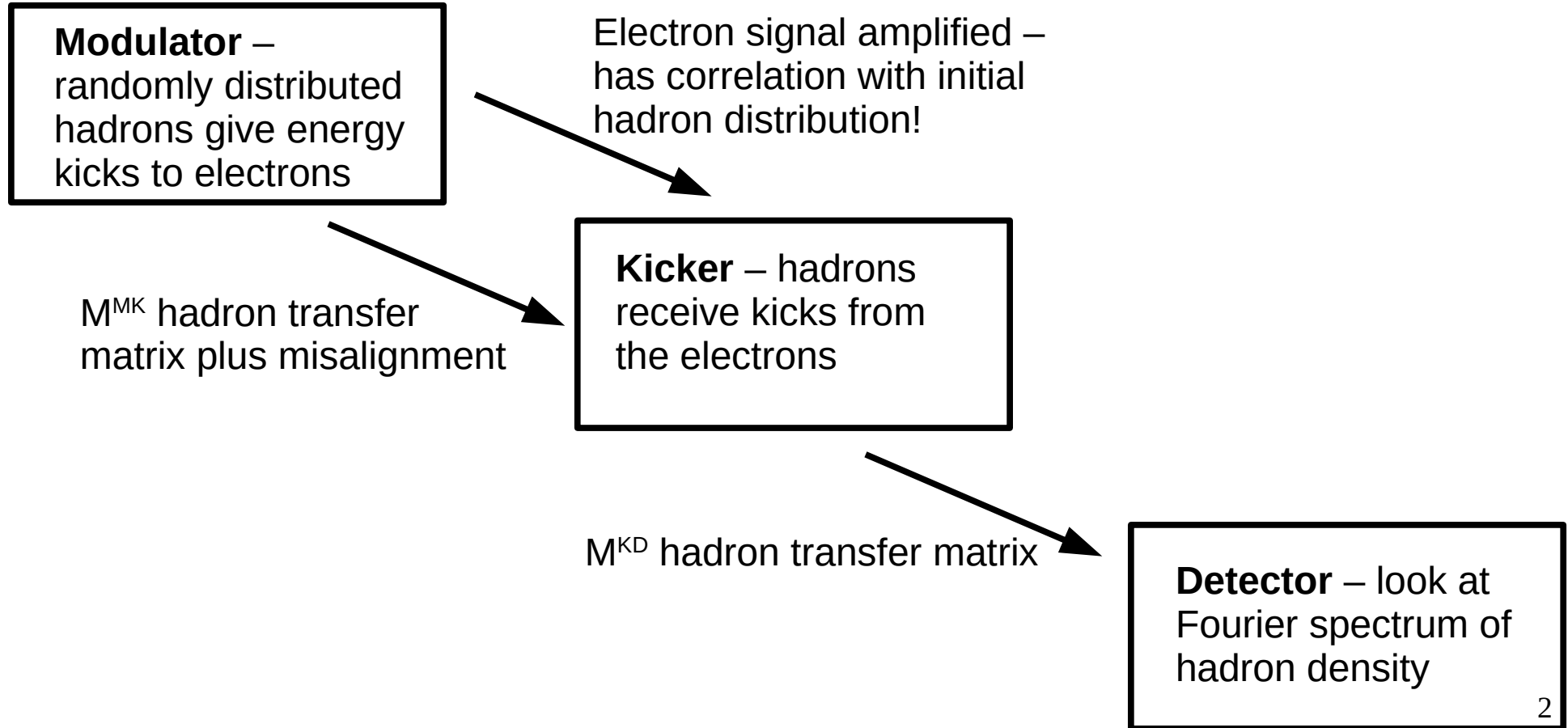


# More Realistic Signal Suppression and Detection

- Full theory, to including transverse dynamics
- Simulation comparisons
- Optimal kicker → detector transfer matrix
- Dipole radiation as a diagnostic?
- Sensitivity to motion within dipole

# Signal Suppression Schematic



# Derivation

- Assuming random shot noise in the modulator, get particle positions in detector
- Obtain densities in Fourier space
- Perform integrals in terms of impedance function
- Plug in analytic expression for MBEC wake function
- Note on notation – repeated “u” indices are summed from 1-6, excluding 5
  - wake defined as fractional momentum kick to hadron

# Derivation (cont.)

$$z_d^{(i)} = z_m^{(i)} + M_{5u}^{MD} \vec{x}_u^{(i)} + M_{56}^{KD} \sum_j w(z_m^{(i)} + M_{5u}^{MK} \vec{x}_u^{(i)} - z_m^{(j)} + \Delta z)$$

$$\rho(z) = \sum_i \delta(z - z_d^{(i)})$$

$$\tilde{\rho}(k) \equiv \int_{-L/2}^{L/2} e^{-ikz} \rho(z) dz = \sum_i e^{-ikz_d^{(i)}}$$

$$|\tilde{\rho}(k)|^2 = \sum_{i,a} e^{-ik[z_d^{(i)} - z_d^{(a)}]}$$

$$|\tilde{\rho}(k)|^2 = N + \sum_{i \neq a} e^{-ik[z_m^{(i)} - z_m^{(a)} + M_{5u}^{MD}(\vec{x}_u^{(i)} - \vec{x}_u^{(a)}) + M_{56}^{KD} \sum_j w(z_m^{(i)} - z_m^{(j)} + M_{5u}^{MK} \vec{x}_u^{(i)} + \Delta z) - M_{56}^{KD} \sum_j w(z_m^{(a)} - z_m^{(j)} + M_{5u}^{MK} \vec{x}_u^{(a)} + \Delta z)]}$$

$$|\tilde{\rho}(k)|^2 = N + \sum_{i \neq a} e^{-ik[z_m^{(i)} - z_m^{(a)} + M_{5u}^{MD}(\vec{x}_u^{(i)} - \vec{x}_u^{(a)})]}$$

$$\times (1 - ikM_{56}^{KD} \sum_j w(z_m^{(i)} - z_m^{(j)} + M_{5u}^{MK} \vec{x}_u^{(i)} + \Delta z) + ikM_{56}^{KD} \sum_j w(z_m^{(a)} - z_m^{(j)} + M_{5u}^{MK} \vec{x}_u^{(a)} + \Delta z))$$

# Derivation (cont.)

$$\begin{aligned}
 \langle |\tilde{\rho}(k)|^2 \rangle &= N + \sum_{i \neq a} \langle e^{-ik[z_{ia} + M_{5u}^{MD}(\vec{x}_u^{(i)} - \vec{x}_u^{(a)})]} \\
 &\times (1 - ikM_{56}^{KD}w(M_{5u}^{MK}\vec{x}_u^{(i)} + \Delta z) + ikM_{56}^{KD}w(M_{5u}^{MK}\vec{x}_u^{(a)} + \Delta z) \\
 &- ikM_{56}^{KD}w(z_{ia} + M_{5u}^{MK}\vec{x}_u^{(i)} + \Delta z) + ikM_{56}^{KD}w(-z_{ia} + M_{5u}^{MK}\vec{x}_u^{(a)} + \Delta z)) \rangle
 \end{aligned}$$

11-D integrals such as (4<sup>th</sup> term):

$$\begin{aligned}
 &-ikM_{56}^{KD} \int_{-L/2}^{L/2} dz_{ia}/L \int_{-\infty}^{\infty} d^5\vec{x}^{(i)} d^5\vec{x}^{(a)} \rho(\vec{x}^{(i)}) \rho(\vec{x}^{(a)}) e^{-ik[z_{ia} + M_{56}^{MD}(\vec{x}_u^{(i)} - \vec{x}_u^{(a)})]} \\
 &\times w(z_{ia} + M_{5u}^{MK}\vec{x}_u^{(i)} + \Delta z)
 \end{aligned}$$

# Derivation (cont.)

$$-ikM_{56}^{KD} \int_{-L/2}^{L/2} dz_{ia}/L \int_{-\infty}^{\infty} d^5\vec{x}^{(i)} d^5\vec{x}^{(a)} \rho(\vec{x}^{(i)}) \rho(\vec{x}^{(a)}) e^{-ik[z_{ia} + M_{56}^{MD}(\vec{x}_u^{(i)} - \vec{x}_u^{(a)})]} \\ \times w(z_{ia} + M_{5u}^{MK} \vec{x}_u^{(i)} + \Delta z)$$

Defining impedance:  $Z(k) \equiv \int_{-L/2}^{L/2} e^{-ikz} w(z) dz$

$$\frac{-ikM_{56}^{KD}}{L} Z(k) e^{ik\Delta z} \int_{-\infty}^{\infty} d^5\vec{x}^{(i)} d^5\vec{x}^{(a)} \rho(\vec{x}^{(i)}) \rho(\vec{x}^{(a)}) e^{-ik[M_{56}^{MD}(\vec{x}_u^{(i)} - \vec{x}_u^{(a)}) - M_{5u}^{MK} \vec{x}_u^{(i)}]}$$

5<sup>th</sup> term (using  $Z(-k) = Z^*(k)$ ):

$$\frac{ikM_{56}^{KD}}{L} Z^*(k) e^{-ik\Delta z} \int_{-\infty}^{\infty} d^5\vec{x}^{(i)} d^5\vec{x}^{(a)} \rho(\vec{x}^{(i)}) \rho(\vec{x}^{(a)}) e^{-ik[M_{56}^{MK} \vec{x}_u^{(a)} + M_{5u}^{MD}(\vec{x}_u^{(i)} - \vec{x}_u^{(a)})]}$$

# Derivation (cont.)

$$\langle |\tilde{\rho}(k)|^2 \rangle = N + 2N \frac{(N-1)}{L} k M_{56}^{KD} [Re(Z(k)) \sin(k\Delta z) + Im(Z(k)) \cos(k\Delta z)] \\ \times e^{-\frac{k^2}{2} \sum_{u \neq 5} \sigma_{\tilde{x}_u}^2 \left[ \left( \tilde{M}_{5u}^{MD} \right)^2 + \left( \tilde{M}_{5u}^{MK} - \tilde{M}_{5u}^{MD} \right)^2 \right]}$$

$$\sigma_{\tilde{x}} \equiv \sqrt{\epsilon_x \beta_x}$$

$$\sigma_{\tilde{x}'} \equiv \sqrt{\epsilon_x / \beta_x}$$

$$\sigma_{\tilde{y}} \equiv \sqrt{\epsilon_y \beta_y}$$

$$\sigma_{\tilde{y}'} \equiv \sqrt{\epsilon_y / \beta_y}$$

$$\sigma_{\tilde{\delta}} \equiv \sigma_{\delta}$$

$$\tilde{M}_{51} \equiv M_{51} - \frac{\alpha_x}{\beta_x} M_{52}$$

$$\tilde{M}_{52} \equiv M_{52}$$

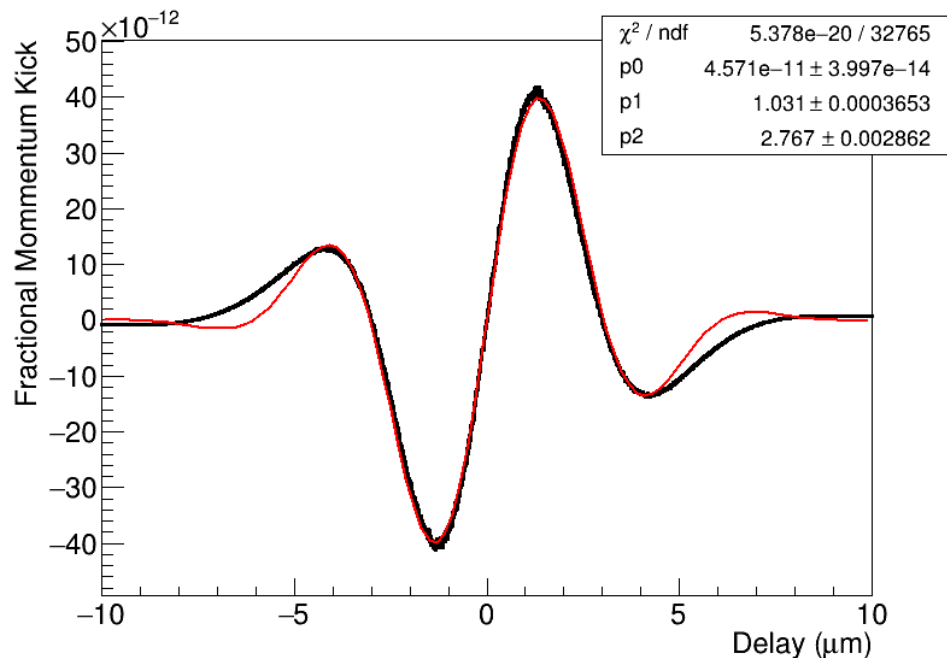
$$\tilde{M}_{53} \equiv M_{53} - \frac{\alpha_y}{\beta_y} M_{54}$$

$$\tilde{M}_{54} \equiv M_{54}$$

$$\tilde{M}_{56} \equiv D_x M_{51} + D'_x M_{52} + D_y M_{53} + D'_y M_{54} + M_{56}$$

$$\frac{\Delta \langle |\tilde{\rho}(k)|^2 \rangle}{\langle |\tilde{\rho}(k)|^2 \rangle} = 2nk M_{56}^{KD} [Re(Z(k)) \sin(k\Delta z) + Im(Z(k)) \cos(k\Delta z)] \\ \times e^{-\frac{k^2}{2} \sum_{u \neq 5} \sigma_{\tilde{x}_u}^2 \left[ \left( \tilde{M}_{5u}^{MD} \right)^2 + \left( \tilde{M}_{5u}^{MK} - \tilde{M}_{5u}^{MD} \right)^2 \right]}$$

# Analytic Wake Model (Saturation On)



$$w(z) = A \sin(\kappa z) e^{-z^2/2\lambda^2}$$

$$A = 4.571 \times 10^{-11}$$

$$\kappa = 1.031/\mu m$$

$$\lambda = 2.767\mu m$$

$$Z(k) = \frac{Ai}{2} \sqrt{2\pi\lambda} \left[ e^{(k+\kappa)^2/2\lambda^2} - e^{(k-\kappa)^2/2\lambda^2} \right]$$

See: S. Nagaitsev et al, “Cooling and diffusion rates in coherent electron cooling concepts”, FERMILAB-CONF-21-054-AD, <https://lss.fnal.gov/archive/2021/conf/fermilab-conf-21-054-ad.pdf> (2021).



# Signal Suppression for Our Wake

$$\frac{\Delta \langle |\tilde{\rho}(k)|^2 \rangle}{\langle |\tilde{\rho}(k)|^2 \rangle} = nk M_{56}^{KD} A \sqrt{2\pi} \lambda \cos(k\Delta z) e^{-\frac{k^2}{2} \sum_{u \neq 5} \sigma_{\tilde{x}_u}^2} \left[ \left( \tilde{M}_{5u}^{MD} \right)^2 + \left( \tilde{M}_{5u}^{MK} - \tilde{M}_{5u}^{MD} \right)^2 \right] \\ \times \left[ e^{-(k+\kappa)^2 \lambda^2 / 2} - e^{-(k-\kappa)^2 \lambda^2 / 2} \right]$$

$$\sigma_{\tilde{x}} \equiv \sqrt{\epsilon_x \beta_x}$$

$$\sigma_{\tilde{x}'} \equiv \sqrt{\epsilon_x / \beta_x}$$

$$\sigma_{\tilde{y}} \equiv \sqrt{\epsilon_y \beta_y}$$

$$\sigma_{\tilde{y}'} \equiv \sqrt{\epsilon_y / \beta_y}$$

$$\sigma_{\tilde{\delta}} \equiv \sigma_{\delta}$$

$$\tilde{M}_{51} \equiv M_{51} - \frac{\alpha_x}{\beta_x} M_{52}$$

$$\tilde{M}_{52} \equiv M_{52}$$

$$\tilde{M}_{53} \equiv M_{53} - \frac{\alpha_y}{\beta_y} M_{54}$$

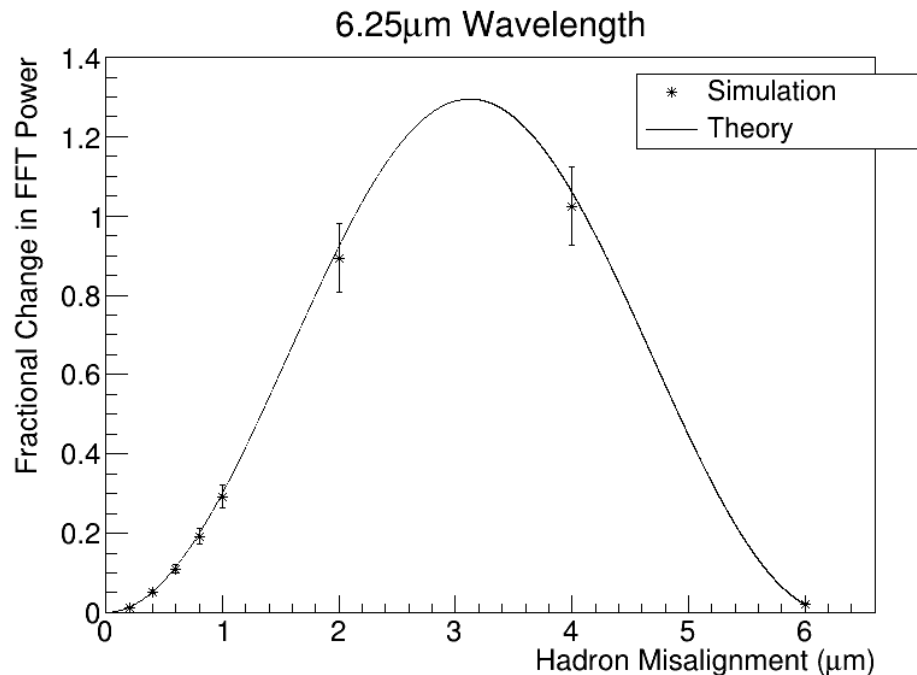
$$\tilde{M}_{54} \equiv M_{54}$$

$$\tilde{M}_{56} \equiv D_x M_{51} + D'_x M_{52} + D_y M_{53} + D'_y M_{54} + M_{56}$$

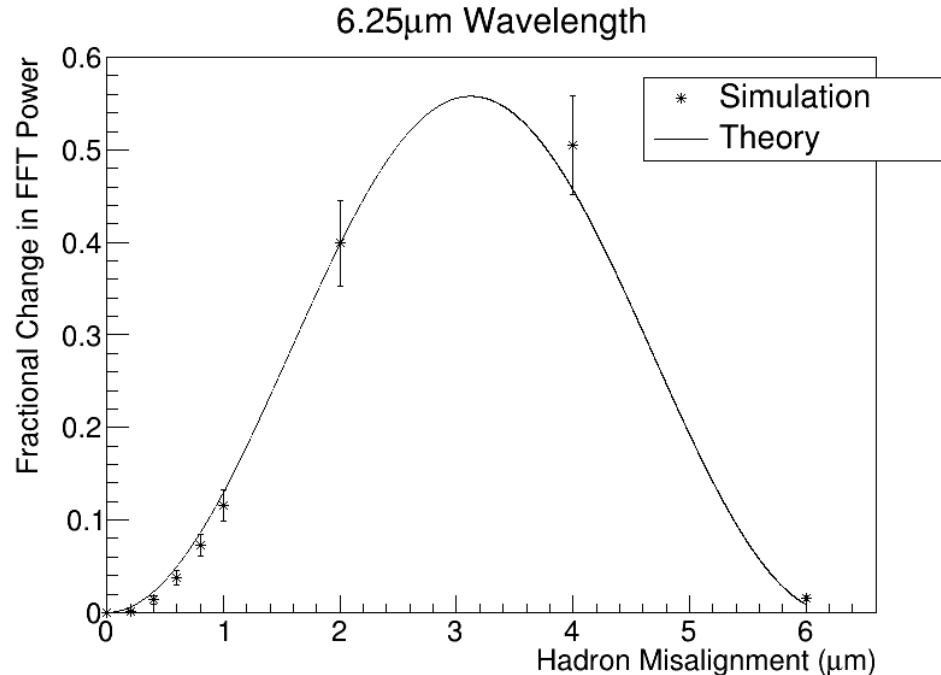
# Comparison with Simulation

( $K \rightarrow D$  Transfer Matrix Inverse of  $M \rightarrow K$ )

(Normalized Relative to No-Offset Case)



Linearized Simulation



Full Simulation  
(Saturation included)

# Optimal Transfer Matrix

- With current modulator → kicker transfer matrix, gradient-descent optimization of signal suppression formula gives maximum (in absolute value sense):

$$\frac{\Delta \langle |\tilde{\rho}(k)|^2 \rangle}{\langle |\tilde{\rho}(k)|^2 \rangle} = -0.22 \cos(k\Delta z)$$

with parameters:

$$M_{51}^{KD} = -8.44 \times 10^{-4} \quad M_{52}^{KD} = -3.00 \times 10^{-2}m \quad M_{56}^{KD} = 2.36 \times 10^{-3}m \quad k = 1.04 \times 10^6/m$$

For comparison, inverse-matrix plots from last slide had used:

$$M_{51}^{KD} = -7.93 \times 10^{-4}, \quad M_{52}^{KD} = -2.86 \times 10^{-2}m, \quad M_{56}^{KD} = 2.26 \times 10^{-3}m, \quad k = 1.06 \times 10^6/m$$

# Detection with Dipole Radiation

# Can We Make Hadron Signal Frequency the Critical Frequency?

- 275GeV protons
- Gamma of 293
- Omega of ck = 3.11e14/s

$$\omega_c = \frac{3\gamma^3 c}{2\rho} \quad (\text{See Jackson E\&M, 3}^{\text{rd}} \text{ ed., eqtn. 14.81})$$

$$B = \frac{E_0}{ec\rho} = \frac{2E_0\omega_c}{3\gamma^3 ec^2} = 25.2T$$

Not realistic!

# Dipole Radiation Estimates

$$\frac{d^2 I}{d\omega d\Omega} = \frac{e^2}{4\pi\epsilon_0 3\pi^2 c} \left(\frac{\omega\rho}{c}\right)^2 \left(\frac{1}{\gamma^2} + \theta^2\right)^2 \left[ K_{2/3}^2(\xi) + \frac{\theta^2}{(1/\gamma^2) + \theta^2} K_{1/3}^2(\xi) \right]$$

$$\xi = \frac{\omega\rho}{3c} \left(\frac{1}{\gamma^2} + \theta^2\right)^{3/2} \quad \text{(See Jackson E\&M, 3<sup>rd</sup> ed., eqtn. 14.76 and 14.79)}$$

- 3.782 T magnets → 243m bending radius (CDR, pg 196)
- Gamma of 293
- Omega of ck = 3.11e14/s
- Look at 0 angle (on-axis)

$$\frac{d^2 I}{d\omega d\Omega} = 1.42 \times 10^{-34} \frac{J}{\frac{1}{s} sr}$$

# Dipole Radiation, Continued

$$\frac{d^2 I}{d\omega d\Omega} = 1.42 \times 10^{-34} \frac{J}{\frac{1}{s} sr}$$

- Assume frequency bandwidth of  $\omega/10$
- Intensity 10m downstream is  $4.5e-23$  J/m<sup>2</sup> per hadron per revolution
- With 7mm/6cm of h<sup>+</sup> bunch overlapping with electrons, core has  $8.1e9$  h<sup>+</sup>
- Get  $3.6e-13$  J/m<sup>2</sup> per bunch per revolution
- With 1160 bunches and revolution time of  $13\mu s$ , have  $3.3e-5$  W/m<sup>2</sup> intensity from core protons

# Black-body Background Radiation

- $3.3\text{e-}5 \text{ W/m}^2$  on-axis radiation intensity from core protons
- If integrate over detector with radius 3.4mm (1/(10 gamma) angular acceptance), a intensity of  **$\sim 3.1\text{e-}5 \text{ W/m}^2$**  10 m downstream – total power of 1.1nW = 35 billion photons/second at detector

$$\frac{dP}{d\omega} = \frac{\hbar}{4\pi^2 c^2} \frac{\omega^3}{e^{\hbar\omega/k_B T} - 1}$$

$$P = \sigma T^4$$

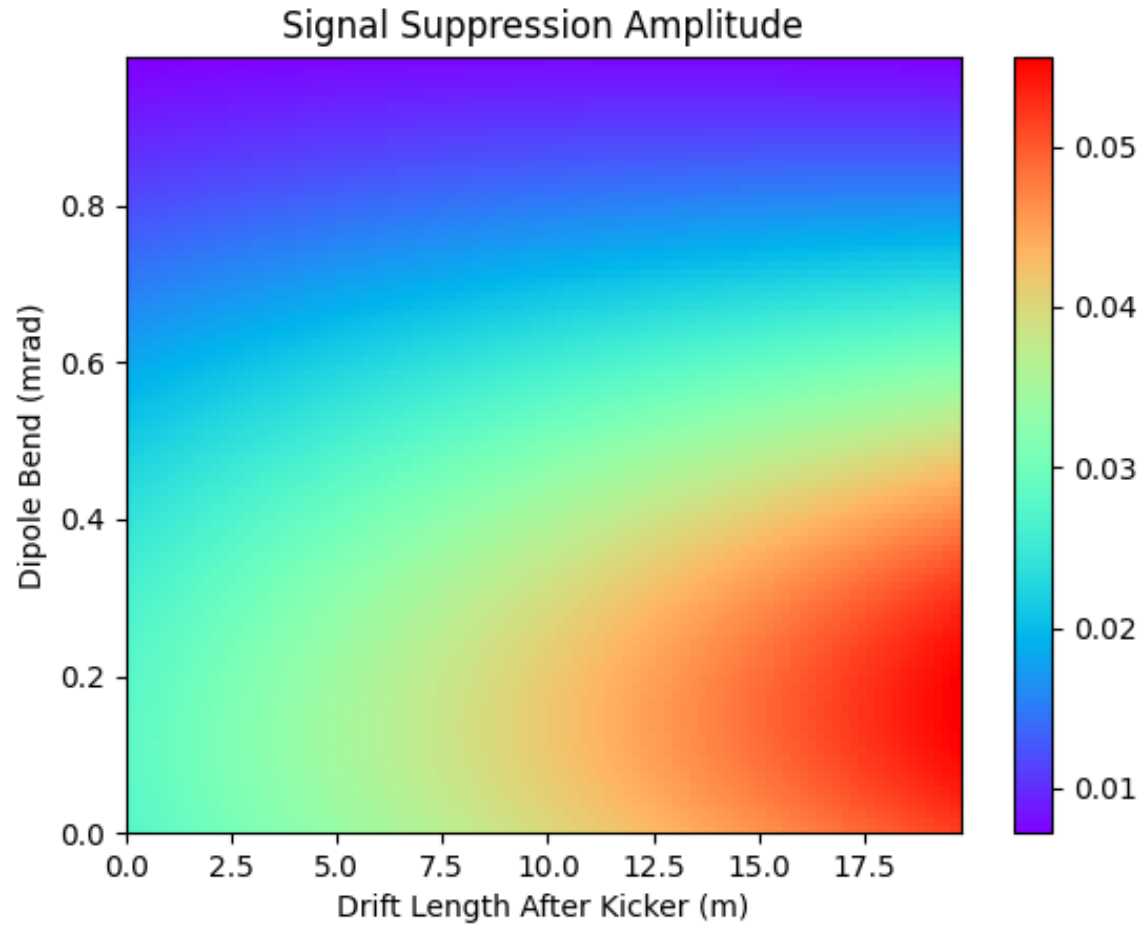
See, eg, J. P. Sethna, “Statistical Mechanics: Entropy, Order Parameters, and Complexity”, eqtn. 7.80 and 7.82

- At 300K, intensity of background in 10% bandwidth is  $10 \text{ W/m}^2$  – swamps any signal
- At 77.29K (LN2), have  $1.7\text{e-}9 \text{ W/m}^2$  within 10% bandwidth – much less than signal  
But, total power over all frequencies is 2 W – again much larger than signal
- At 4.22K (LHe), total power over all frequencies is  $1.8\text{e-}5 \text{ W}$  – comparable to signal



# Drift and Dipole to Generate R56?

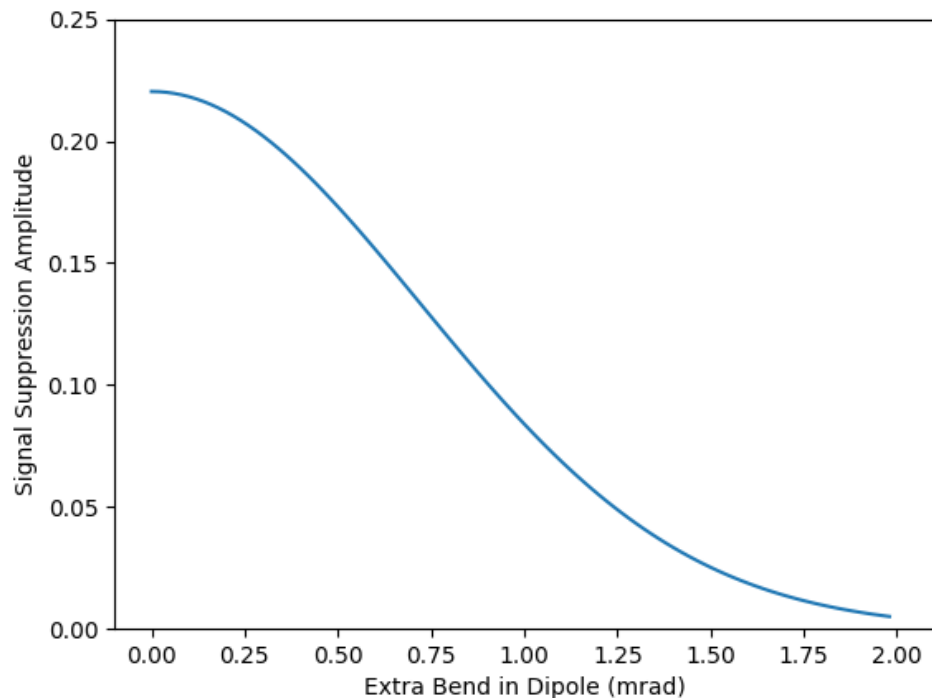
- Need to also generate proper transfer matrix from kicker to detector
- What can be achieved by a drift and dipole?
- Use dipole of strength given previously, scan bend angle within dipole and drift length before (plus  $L_K/2$  unavoidable drift from kicker)
- Ignore edge focusing of dipole
- Amplitudes of signal suppression defined by  $-A \cos(k\Delta z)$



Need large drift, and little bending

Note  $\frac{1}{10\gamma} \approx 0.3 \text{ mrad}$

# Sensitivity to Extra Bend in Dipole (Optimal Transfer Matrix)



Note:  $\frac{1}{10\gamma} \approx 0.3mrad$

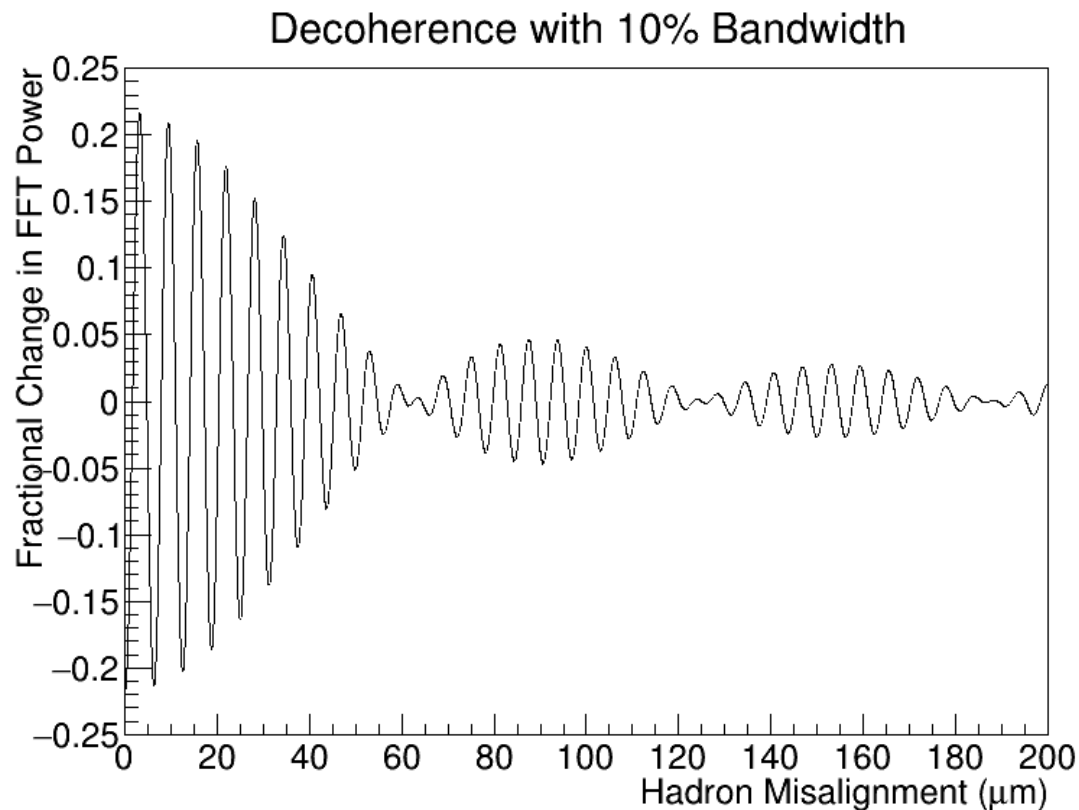
# Conclusions

- Can get optimal signal suppression amplitude of  $\sim 22\%$
- Can get  $\sim 1/10$  of optimal power with only a short drift and dipole
- Operation at critical frequency not possible with a pure dipole
- Cryogenic detector likely needed so that we aren't swamped by thermal noise
- Signal output is sensitive to where we observe in dipole, on scale comparable to radiation opening angle, leading to up to a few percent reduction in signal

# Backup Slides

- Decoherence of signal suppression
- Undulator radiation
- Parameter table

# Decoherence of Signal Suppression



$$\langle \cos(k\Delta z) \rangle = \frac{1}{\Delta k} \int_k^{k+\Delta k} \cos(k'\Delta z) dk' = \frac{\sin((k + \Delta k)\Delta z) - \sin(k\Delta z)}{\Delta k \Delta z}$$

# Undulator Radiation?

$$K = \frac{eB\lambda_u}{2\pi mc}$$

$$\lambda = \frac{\lambda_u}{2\gamma^2} (1 + K^2/2)$$

- For simplicity, say undulator field is also 3.782T  
Then,  $K = 0.20$ , need undulator period of 1.02m

# Undulator Radiation?

$$\frac{dP}{d\omega\Omega} = \hbar\omega\alpha N^2\gamma^2 \frac{1}{\omega} \frac{I}{e} \frac{K^2}{(1 + K^2/2)^2} \left[ J_0\left(\frac{K^2}{4(1 + K^2/2)}\right) - J_1\left(\frac{K^2}{4(1 + K^2/2)}\right) \right]^2$$

(Assuming operation at first harmonic – see lecture 12 of [https://www.slac.stanford.edu/~xiahuang/USPAS\\_2017.htm](https://www.slac.stanford.edu/~xiahuang/USPAS_2017.htm))

- For N=1 undulator period, 10% bandwidth, detector 10m from undulator, we get  $\sim 5.5 \times 10^{-4} \text{ W/m}^2$  intensity on-axis from core protons  
(cf.  $3.3 \times 10^{-5} \text{ W/m}^2$  for dipole radiation)
- Does not change any of our conclusions about thermal noise – still lose to  $10 \text{ W/m}^2$  thermal noise at 300K unless we have  $\sim 100$  periods (impossible in arcs)
- Undulator likely not worth the extra effort (not considered: beam evolution within undulator also makes this tricky, as seen in dipole case)



Case	100 GeV	275 GeV
<b><i>Geometry</i></b>		
Modulator Length (m)	45.420181	45.420181
Kicker Length (m)	45.420181	45.420181
Number of Amplifier Drifts	2	2
Amplifier Drift Lengths (m)	36.579819	36.579819
<b><i>Proton Parameters</i></b>		
Protons per Bunch	6.9e10	6.9e10
Proton Bunch Length (cm)	7	6
Proton Fractional Energy Spread	9.7e-4	6.8e-4
Proton Emittance (x/y) (nm)	30 / 2.7	11.3 / 1
Horizontal/Vertical Proton Betas in Modulator (m)	20.96904 / 55.920402	39 / 39
Horizontal/Vertical Proton Dispersion in Modulator (m)	0.848258 / 0.235044	1 / 0
Horizontal/Vertical Proton Dispersion Derivative in Modulator (m)	-0.023971 / 0	-0.023 / 0
Horizontal/Vertical Proton Betas in Kicker (m)	20.96904 / 55.920402	39 / 39
Horizontal/Vertical Proton Dispersion in Kicker (m)	0.848258 / 0.235044	1 / 0
Horizontal/Vertical Proton Dispersion Derivative in Kicker (m)	0.023971 / 0	0.023 / 0
Proton Horizontal/Vertical Phase Advance (rad)	5.145 / 4.9	4.7895 / 4.7895
R56 in Proton Chicane (mm)	-6	-2.25956
<b><i>Electron Parameters</i></b>		
Electron Bunch Charge (nC)	1	1
Electron Bunch Length (mm)	15	7
Electron Peak Current (A)	~ 8	~ 17
Electron Fractional Slice Energy Spread	1e-4	1e-4
Electron Normalized Emittance (x/y) (mm-mrad)	2.8 / 2.8	2.8 / 2.8
Horizontal/Vertical Electron Betas in Modulator (m)	60.043513 / 84.017684	40 / 20
Horizontal/Vertical Electron Betas in Kicker (m)	10.006454 / 13.977728	4 / 4
Horizontal/ Vertical Electron Betas in Amplifiers (m)	1.537031 / 1.537031	1 / 1
R56 in First Electron Chicane (mm)	15.786	5.0
R56 in Second Electron Chicane (mm)	15.786	5.0
R56 in Third Electron Chicane (mm)	-42.573	-11.45
<b><i>Cooling Times</i></b>		
Horizontal/Vertical/Longitudinal IBS Times (hours)	2.0 / 4.0 / 2.5	2.0 / 5? / 2.9
Horizontal/Vertical/Longitudinal Cooling Times (hours)	1.7 / 3.8 / 2.2	1.8 / $\infty$ / 2.8